

# Some discrete integrable equations related to an elliptic curve

V.E. Adler, Yu.B. Suris

13th May 2005

## Contents

1	Introduction	2
2	Nonlinear superposition for the Krichever-Novikov equation	4
3	Three-leg form of $(Q_4)$	7
4	$S_4$ symmetry group of $(Q_4)$	8
5	Three-leg form implies $3D$ -consistency	9
6	$3D$ -consistency $\rightarrow$ zero curvature representation	11
7	Discrete Toda lattices	12
8	Elliptic Toda lattice	20
9	Elliptic Ruijsenaars-Toda lattice	22

---

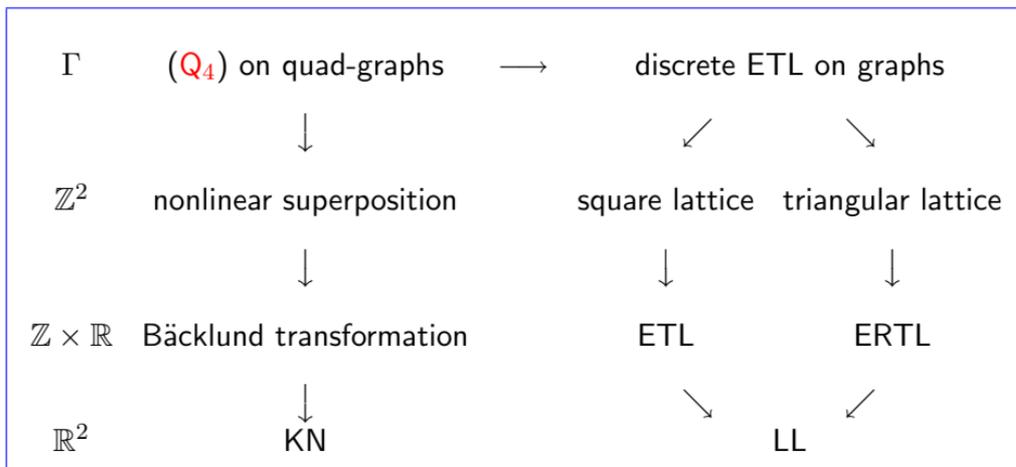
[1] V.E. Adler, Yu.B. Suris.  $Q_4$ : Integrable Master Equation Related to an Elliptic Curve. *Int. Math. Res. Not.* **47** (2004) 2523–2553.

# 1 Introduction

We discuss several 2D integrable systems with the spectral parameter on an elliptic curve: Landau-Lifshitz (LL) and Krichever-Novikov (KN) equations, elliptic Toda lattice (ETL) and elliptic Ruijsenaars-Toda lattice (ERTL). These models can be unified on the basis of a single discrete equation (subscripts denote the shifts in the square lattice):

$$(Q_4) \quad abc(uu_1u_2u_{12} + 1) + a(uu_1 + u_2u_{12}) - b(uu_2 + u_1u_{12}) - c(uu_{12} + u_1u_2) = 0$$

$$A^2 = a^4 - da^2 + 1, \quad B^2 = b^4 - db^2 + 1, \quad c = \frac{aB - Ab}{1 - a^2b^2}.$$



**Remark.** Eq. (Q<sub>4</sub>) was introduced in [2] and studied further in [3, 4, 1]. However, in these papers it was presented in much more cumbersome form related to the Weierstrass form of the elliptic curve ( $A^2 = 4a^3 - g_2a - g_3$ ). The given form of equation (Q4) was found by Hietarinta [SIDE-2004 talk].

- 
- [2] V.E. Adler. Bäcklund transformation for the Krichever-Novikov equation. *Int. Math. Res. Notices* **1998:1** 1–4.
- [3] V.E. Adler, A.I. Bobenko, Yu.B. Suris. Classification of integrable equations on quad-graphs. The consistency approach. *Comm. Math. Phys.* **233** (2003) 513–543.
- [4] F. Nijhoff. Lax pair for the Adler (lattice Krichever-Novikov) system. *Phys. Lett. A* **297** (2002) 49–58.

## 2 Nonlinear superposition for the Krichever-Novikov equation

The Krichever-Novikov equation [5]

$$(1) \quad u_t = u_{xxx} - \frac{3(u_{xx}^2 - r(u))}{2u_x}, \quad r^{(5)} = 0$$

is the most generic nonlinear integrable equation of the form  $u_t = u_{xxx} + f(u_{xx}, u_x, u)$ . Accordingly to [6], all other equations of this type are related via differential substitutions to eq (1) when some zeroes of the polynomial  $r$  are multiple, while the case of simple zeroes is isolated.

Bäcklund auto-transformation for (1) is of the form [2]

$$(2) \quad u_x \tilde{u}_x = h(u, \tilde{u})$$

where  $h$  is the biquadratic polynomial in  $u, \tilde{u}$ , such that

$$r(u) = h_{\tilde{u}}^2 - 2hh_{\tilde{u}\tilde{u}}, \quad r(\tilde{u}) = h_u^2 - 2hh_{uu}.$$

---

[5] I.M. Krichever, S.P. Novikov. Holomorphic bundles over algebraic curves and nonlinear equations. *Uspekhi Mat. Nauk* **35:6** (1980) 47–68.

[6] S.I. Svinolupov, V.V. Sokolov, R.I. Yamilov. Bäcklund transformations for integrable evolution equations. *Dokl. Akad. Nauk SSSR* **271:4** (1983) 802–805.

The polynomial  $h$  corresponding to  $r(u) = u^4 - du^2 + 1$  depends on an additional parameter  $(a, A)$  on the elliptic curve  $A^2 = r(a)$ :

$$h(u, \tilde{u}; a, A) = \frac{1}{2a}(a^2u^2\tilde{u}^2 - 2Au\tilde{u} - u^2 - \tilde{u}^2 + a^2).$$

BTs corresponding to the different values of  $\alpha$  commute, and eq. (Q<sub>4</sub>) defines the nonlinear superposition principle for these BT. This means the following. The equations

$$(3) \quad \begin{aligned} u_x u_{1,x} &= h(u, u_1; a, A) & u_{2,x} u_{12,x} &= h(u_2, u_{12}; a, A) \\ u_x u_{2,x} &= h(u, u_2; b, B) & u_{1,x} u_{12,x} &= h(u_1, u_{12}; b, B) \end{aligned}$$

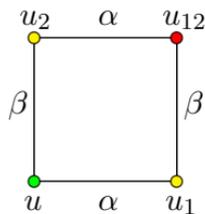
imply the reducible constraint

$$h(u, u_1; a, A)h(u_2, u_{12}; a, A) - h(u, u_2; b, B)h(u_1, u_{12}; b, B) = Q\tilde{Q} = 0$$

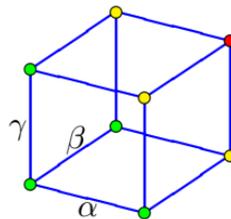
where  $Q$  is the l.h.s. of the eq. (Q<sub>4</sub>) and  $\tilde{Q} = Q|_{b \rightarrow -b}$ . The statement is that the constraint  $Q = 0$  is consistent with the dynamics on  $x$  defined by (3):  $\frac{dQ}{dx} \Big|_{Q=0} = 0$ .

The important property of  $(Q_4)$  eq. is  $3D$ -consistency, or consistency around a cube. This property means that if we assign 6 copies of  $(Q_4)$  to the faces of a cube, with the common values of the parameters on the edges, then for arbitrary choice of initial data  $u, u_1, u_2, u_3$  the values  $u_{123}$  calculated in three possible ways coincide.

The classification of  $3D$ -consistent equation, under some additional assumptions, was obtained in [9].



nonlinear superposition

 $3D$ -consistency

- 
- [7] F.W. Nijhoff, A.J. Walker. The discrete and continuous Painlevé hierarchy and the Garnier system. *Glasgow Math. J.* **43A** (2001) 109–123.
- [8] A.I. Bobenko, Yu.B. Suris. Integrable systems on quad-graphs. *Int. Math. Res. Notices* **11** (2002) 573–611.
- [9] V.E. Adler, A.I. Bobenko, Yu.B. Suris. Classification of integrable equations on quad-graphs. The consistency approach. *Comm. Math. Phys.* **233** (2003) 513–543.

### 3 Three-leg form of (Q<sub>4</sub>)

**Statement 1.** Equation (Q<sub>4</sub>) is equivalent, under the changes

$$a = \delta^{-1} \operatorname{sn} \alpha, \quad A = \operatorname{sn}' \alpha, \quad b = \delta^{-1} \operatorname{sn} \beta, \quad B = \operatorname{sn}' \beta, \quad c = \delta^{-1} \operatorname{sn}(\alpha - \beta), \quad u = \delta^{-1} \operatorname{sn} q$$

(where  $\operatorname{sn} x \equiv \operatorname{sn}(x, d^{-2})$ ), to the equation

$$(4) \quad F(q, q_1, \alpha) / F(q, q_2, \beta) = F(q, q_{12}, \alpha - \beta)$$

where

$$F(q, \tilde{q}, \alpha) = \frac{\operatorname{sn}(q + \alpha) - \operatorname{sn}(\tilde{q})}{\operatorname{sn}(q - \alpha) - \operatorname{sn}(\tilde{q})} \cdot \frac{\Theta_4(q + \alpha)}{\Theta_4(q - \alpha)}$$

This property implies several important consequences:  $S_4$  symmetry group of the equation,  $3D$ -consistency, relation to discrete Toda type lattices.

## 4 $S_4$ symmetry group of $(Q_4)$

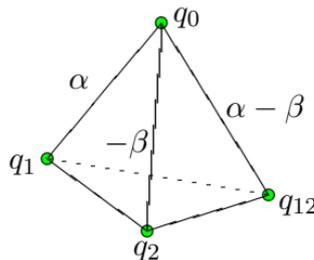
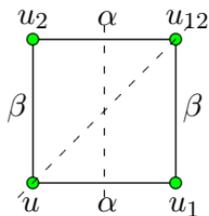
Obviously, eq  $(Q_4)$  admits the symmetry group  $D_4$  of the square:

$$Q(u, u_1, u_2, u_{12}, \alpha, \beta) = Q(u_1, u, u_{12}, u_2, \alpha, \beta) = -Q(u, u_2, u_1, u_{12}, \beta, \alpha).$$

Due to this symmetry the three-leg form can be centered in an arbitrary vertex. On the other hand, three-leg form exhibits one more symmetry which is hidden in the rational form: the diagonals of the quadrilateral are on the equal footing with its edges.

Indeed, due to the property  $F(q, \tilde{q}, \alpha) = 1/F(q, \tilde{q}, -\alpha)$ , eq  $(4)$  can be rewritten in the form

$$F(q, q_1, \alpha)F(q, q_2, -\beta)F(q, q_{12}, \beta - \alpha) = 1.$$



## 5 Three-leg form implies 3D-consistency

**Theorem 2.** Eq (Q<sub>4</sub>) is 3D-consistent.

*Proof.* Consider the system of equations, associated to the faces of the cube:

$$\begin{array}{ll}
 (E_{12}) & Q(u, u_1, u_2, u_{12}, \alpha_1, \alpha_2) = 0 \\
 (E_{13}) & Q(u, u_1, u_3, u_{13}, \alpha_1, \alpha_3) = 0 \\
 (E_{23}) & Q(u, u_2, u_3, u_{23}, \alpha_2, \alpha_3) = 0
 \end{array}
 \qquad
 \begin{array}{ll}
 (\tilde{E}_{12}) & Q(u_3, u_{13}, u_{23}, u_{123}, \alpha_1, \alpha_2) = 0 \\
 (\tilde{E}_{13}) & Q(u_2, u_{13}, u_{23}, u_{123}, \alpha_1, \alpha_3) = 0 \\
 (\tilde{E}_{23}) & Q(u_1, u_{12}, u_{13}, u_{123}, \alpha_2, \alpha_3) = 0
 \end{array}$$

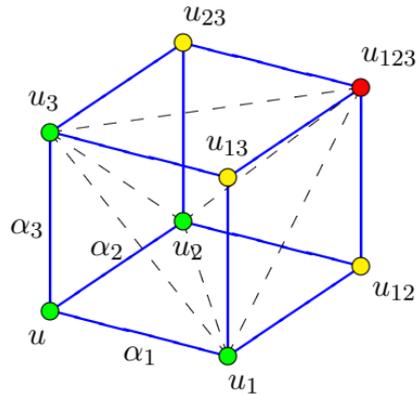
One have to prove that if the values  $u_{12}, u_{13}, u_{23}$  are defined from eqs on the left, for arbitrary initial data  $u, u_1, u_2, u_3$ , then the rest eqs define one and the same value  $u_{123}$ .

It is enough to show that if  $u_{123}$  is defined by eq ( $\tilde{E}_{23}$ ), then ( $\tilde{E}_{13}$ ) is fulfilled as well. Rewrite eqs, containing  $u_1$ , in three-leg forms:

$$\begin{aligned}
 F(q_1, q_{13}, \alpha_3)/F(q_1, q, \alpha_1) &= F(q_1, q_3, \alpha_3 - \alpha_1), \\
 F(q_1, q_{12}, \alpha_2)/F(q_1, q, \alpha_1) &= F(q_1, q_2, \alpha_2 - \alpha_1), \\
 F(q_1, q_{12}, \alpha_2)/F(q_1, q_{13}, \alpha_3) &= F(q_1, q_{123}, \alpha_2 - \alpha_3).
 \end{aligned}$$

From here the equation

$$F(q_1, q_2, \alpha_2 - \alpha_1)/F(q_1, q_3, \alpha_3 - \alpha_1) = F(q_1, q_{123}, \alpha_2 - \alpha_3)$$



follows, relating the fields at the vertices of the dashed tetrahedron. This is nothing but the three-leg form of the equation

$$Q(u_1, u_2, u_3, u_{123}, \alpha_2 - \alpha_1, \alpha_2 - \alpha_3) = 0,$$

centered at  $q_1$ . This can be centered at  $q_2$ , as well, resulting in the cyclic shift of indices:

$$F(q_2, q_3, \alpha_3 - \alpha_2)/F(q_2, q_1, \alpha_1 - \alpha_2) = F(q_2, q_{123}, \alpha_3 - \alpha_1).$$

The latter equation, together with the three-leg forms of equations  $(E_{12})$ ,  $(E_{23})$  centered at  $q_2$ , leads to the three-leg form of  $(\tilde{E}_{13})$ , as required.  $\square$

## 6 $3D$ -consistency $\rightarrow$ zero curvature representation

An affine-linear equation  $Q = 0$  may be interpreted as Möbius transformation between any pair of variables, with coefficients depending on the rest pair. Let

$$u_{13} = M(u_1, u, \alpha_1, \alpha_3)[u_3] = \frac{Au_3 + B}{Cu_3 + D}$$

then

$$u_{23} = M(u_2, u, \alpha_2, \alpha_3)[u_3], \quad u_{123} = M(u_{12}, u_2, \alpha_1, \alpha_3)[u_{23}] = M(u_{12}, u_1, \alpha_2, \alpha_3)[u_{13}].$$

Since the composition of Möbius transformations corresponds to the product of the matrices, hence denoting  $\alpha_3 = \lambda$  and introducing the normalization factor yields the zero curvature representation

$$L(u_{12}, u_1, \alpha_2, \lambda)L(u_1, u, \alpha_1, \lambda) = L(u_{12}, u_2, \alpha_1, \lambda)L(u_2, u, \alpha_2, \lambda)$$

with the matrix

$$L(u_1, u, \alpha_1, \lambda) = (AD - BC)^{-1/2} \begin{pmatrix} A & B \\ C & D \end{pmatrix}.$$

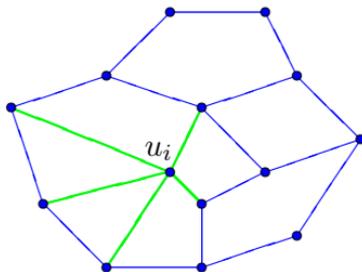
For the  $(Q_4)$  equation one obtains ( $a = d^{-1} \operatorname{sn} \alpha$ ,  $l = d^{-1} \operatorname{sn} \lambda$ ,  $m = d^{-1} \operatorname{sn}(\alpha - \lambda)$ )

$$L(u_1, u, \alpha_1, \lambda) = h(u, u_1, \alpha)^{-1/2} \begin{pmatrix} lu + mu_1 & -alm - auu_1 \\ almuu_1 + a & -lu_1 - mu \end{pmatrix}.$$

## 7 Discrete Toda lattices

Discrete Toda lattices can be defined as equations on “stars” for arbitrary planar graph  $G$ :

$$\sum_{j:(i,j) \in E_G} f_{ij}(u_i, u_j) = 0$$



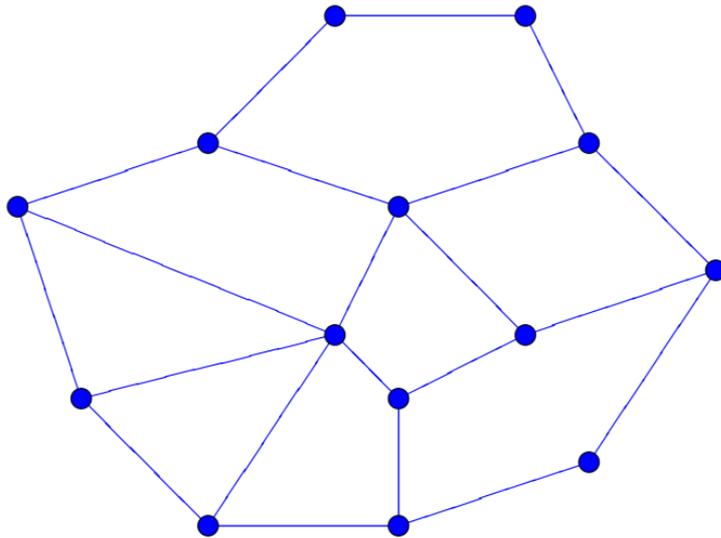
There exists the correspondence between the planar graphs  $G$  and the bipartite quad-graphs  $Q$ :

$$V_Q = V_G \cup V_{G^*}, \quad E_Q = \{(i, i^*) \mid i \in V_G, i^* \in V_{G^*}, i \in f(i^*)\}$$

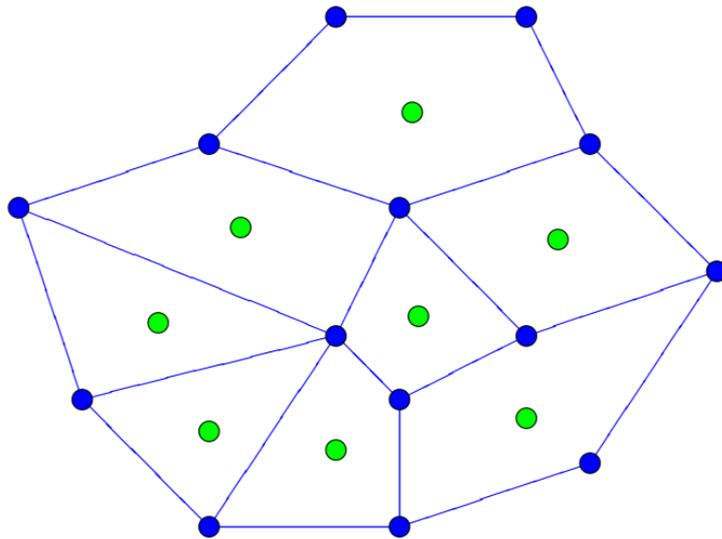
where  $f(i^*)$  is the face of  $G$ , corresponding to the vertex  $i^*$  of dual graph. In other words, the edges of  $G$  are diagonals of the faces of  $Q$  joining the vertices of one of two types.

—  $E_G$

●  $V_G$



- $E_G$
- $V_G$
- $V_{G^*}$

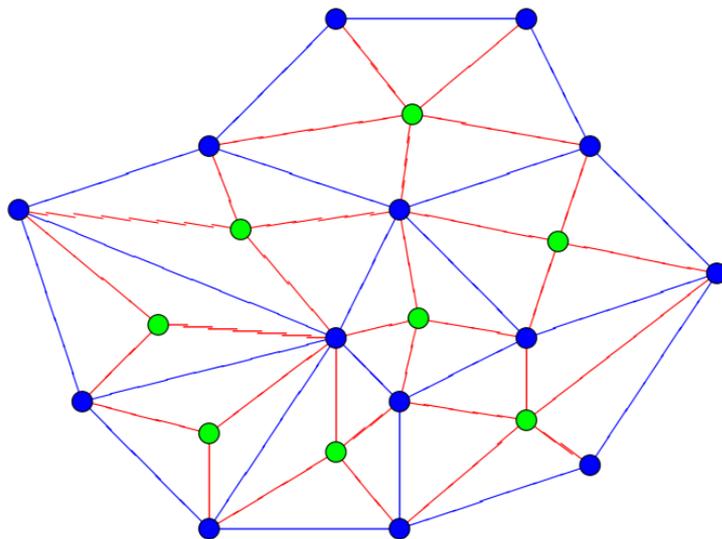


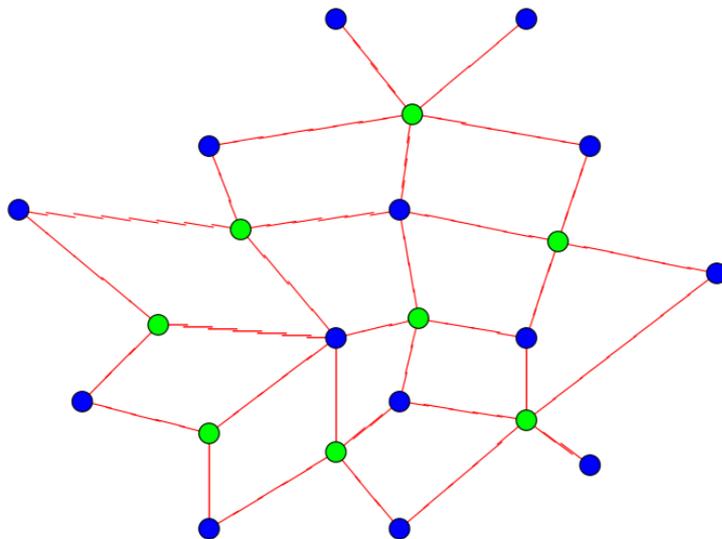
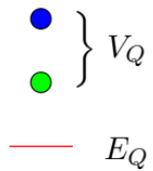
—  $E_G$

●  $V_G$

●  $V_{G^*}$

—  $E_Q$

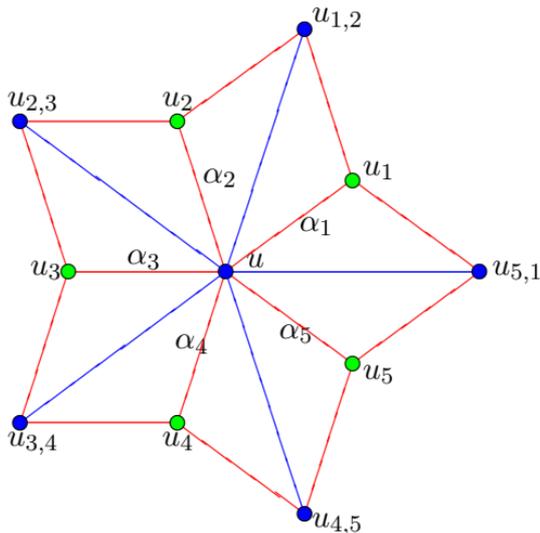




Now, associate eq  $(Q_4)$  to each face of a quad-graph. In all faces, incident to some blue vertex, consider the three-leg form of equation centered in this vertex. Then the product of these equations is free from the variables associated to the green vertices:

$$(5) \quad \prod_{j=1}^n F(q, q_{j,j+1}, \alpha_{j+1} - \alpha_j) = 1$$

In the original variables  $u$  this equation is **rational**.



The eliminated variables  $u_j$  play the role of  $\psi$ -functions: their telescopic cancellation means that the composition of the Möbius transformations

$$u_2 = M(u, u_{1,2}, \alpha_1, \alpha_2)[u_1], \quad u_3 = M(u, u_{2,3}, \alpha_2, \alpha_3)[u_2], \quad \dots$$

$$u_n = M(u, u_{n-1,n}, \alpha_{n-1}, \alpha_n)[u_{n-1}], \quad u_1 = M(u, u_{n,1}, \alpha_n, \alpha_1)[u_n]$$

is the identity transformation. In the matrix language,

$$\prod_{1 \leq j \leq n}^{\curvearrowright} L(u, u_{j,j+1}, \alpha_j + \lambda, \alpha_{j+1} + \lambda) = \text{const } I$$

where  $\lambda$  appears due to the shift invariance of (5).

**Remark.** The 3D-consistent equations were classified in [3], under some additional assumptions. The legs for the other equations of the list can be obtained as limiting cases:

	$F(q, \tilde{q}, \alpha)$	$u = u(q)$	$a = a(\alpha)$
$(Q_1)_{\delta=0}$	$\exp(\alpha/(q - \tilde{q}))$	$q$	$\alpha$
$(Q_1)_{\delta=1}$	$\frac{q - \tilde{q} + \alpha}{q - \tilde{q} - \alpha}$	$q$	$\alpha$
$(Q_2)$	$\frac{(q + \tilde{q} + \alpha)(q - \tilde{q} + \alpha)}{(q + \tilde{q} - \alpha)(q - \tilde{q} - \alpha)}$	$q^2$	$\alpha$
$(Q_3)_{\delta=0}$	$\frac{\sinh(q - \tilde{q} + \alpha)}{\sinh(q - \tilde{q} - \alpha)}$	$\exp 2q$	$\exp 2\alpha$
$(Q_3)_{\delta=1}$	$\frac{\sinh(q + \tilde{q} + \alpha) \sinh(q - \tilde{q} + \alpha)}{\sinh(q + \tilde{q} - \alpha) \sinh(q - \tilde{q} - \alpha)}$	$\cosh 2q$	$\exp 2\alpha$

The leg for the version of  $(Q_4)$  corresponding to the elliptic curve in Weierstrass form:

$$(6) \quad F(q, \tilde{q}, \alpha) = \frac{\sigma(q + \tilde{q} + \alpha)\sigma(q - \tilde{q} + \alpha)}{\sigma(q + \tilde{q} - \alpha)\sigma(q - \tilde{q} - \alpha)}.$$

## 8 Elliptic Toda lattice

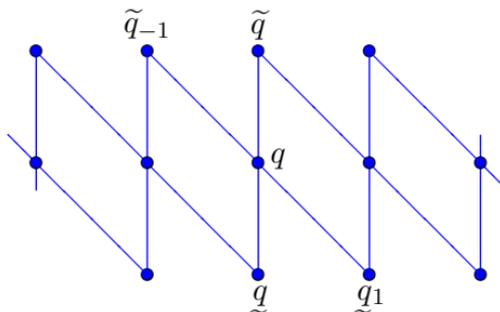
Consider the discrete Toda system on the skew square lattice:

$$F(q, \tilde{q}, \varepsilon)F(q, \tilde{q}_{-1}, -\varepsilon)F(q, \underline{q}, \varepsilon)F(q, \underline{q}_1, -\varepsilon) = 1$$

([10], in the rational variables  $u$ ). It can be written in the Hamiltonian form

$$p = -f(q, \tilde{q}, \varepsilon) + f(q, \tilde{q}_{-1}, \varepsilon), \quad \tilde{p} = f(\tilde{q}, q, \varepsilon) - f(\tilde{q}, \underline{q}_1, \varepsilon)$$

where  $f = \log F$ . Here and in the next section we will use the leg in the form (6).



[10] V.E. Adler. Discretizations of the Landau-Lifshitz equation. *Teor. Math. Phys.* **124:1** (2000) 897–908.

Consider the continuous limit  $\tilde{q} = q + \varepsilon q_x$ . Taking into account the relations

$$\lim_{\varepsilon \rightarrow 0} f(\tilde{q}, q, \varepsilon) = \log \frac{q_x + 1}{q_x - 1}, \quad \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (f(\tilde{q}, q, \varepsilon) + f(q, \tilde{q}, \varepsilon)) = 2\zeta(2q),$$

one finds

$$p = \log \frac{q_x + 1}{q_x - 1},$$

$$p_x = -\zeta(q + q_1) + \zeta(q_1 - q) - \zeta(q + q_{-1}) - \zeta(q - q_{-1}) + 2\zeta(2q).$$

From here the Newtonian equations follow:

$$\frac{q_{xx}}{q_x^2 - 1} = \zeta(q_1 + q) - \zeta(q_1 - q) + \zeta(q + q_{-1}) + \zeta(q - q_{-1}) - 2\zeta(2q).$$

This is the [elliptic Toda lattice](#) as given in [11]. The rational form of this equation [12, 13] reads

$$\frac{u_{xx} - r'(u)/2}{u_x^2 - r(u)} = \frac{1}{u - u_1} + \frac{1}{u - u_{-1}}.$$

[11] I.M. Krichever. Elliptic analog of the Toda lattice. *Int. Math. Res. Notices* **8** (2000) 383–412.

[12] A.B. Shabat, R.I. Yamilov. Symmetries of nonlinear chains, *Len. Math. J.* **2:2** (1991) 377.

[13] R.I. Yamilov. Classification of Toda type scalar lattices, Proc. NEEDS'93, World Scientific Publ., Singapore, 1993, 423–431.

## 9 Elliptic Ruijsenaars-Toda lattice

Analogously, the Hamiltonian form of the discrete Toda system on the triangular lattice is

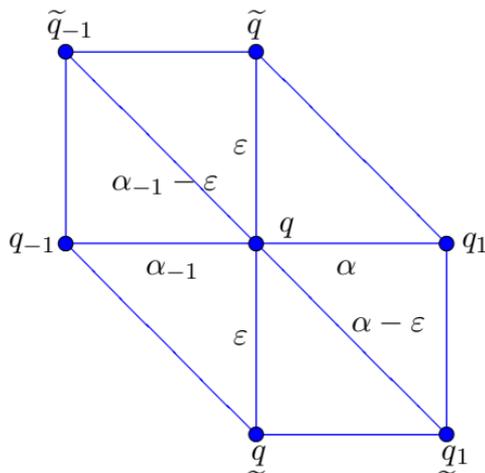
$$p = -f(q, \tilde{q}, \varepsilon) + f(q, q_1, \alpha) + f(q, q_{-1}, \alpha_{-1}) - f(q, \tilde{q}_{-1}, \alpha_{-1} - \varepsilon),$$

$$\tilde{p} = f(\tilde{q}, q, \varepsilon) + f(\tilde{q}, q_1, \alpha - \varepsilon).$$

Under the continuous limit one obtains

$$p = \log \frac{q_x + 1}{q_x - 1} + f(q, q_1, \alpha),$$

$$p_x = q_x f_q(q, q_1, \alpha) + q_{-1,x} f_{q_{-1}}(q, q_{-1}, \alpha_{-1}) - f_\alpha(q, q_1, \alpha) - f_{\alpha_{-1}}(q, q_{-1}, \alpha_{-1}) + 2\zeta(2q).$$



From here the Newtonian equations of the [elliptic Ruijsenaars-Toda lattice](#) follow:

$$\frac{q_{xx}}{q_x^2 - 1} = q_{1,x}f_{q_1}(q, q_1, \alpha) - q_{-1,x}f_{q_{-1}}(q, q_{-1}, \alpha_{-1}) + f_{\alpha}(q, q_1, \alpha) + f_{\alpha_{-1}}(q, q_{-1}, \alpha_{-1}) - 2\zeta(2q)$$

where

$$\begin{aligned} 2f_{q_1}(q, q_1, \alpha) &= \zeta(q + q_1 + \alpha) - \zeta(q - q_1 + \alpha) - \zeta(q + q_1 - \alpha) + \zeta(q - q_1 - \alpha), \\ 2f_{\alpha}(q, q_1, \alpha) &= \zeta(q + q_1 + \alpha) + \zeta(q - q_1 + \alpha) + \zeta(q + q_1 - \alpha) + \zeta(q - q_1 - \alpha). \end{aligned}$$

The rational form of this lattice [14] reads

$$\frac{2u_{xx} - r'(u)}{u_x^2 - r(u)} = -\frac{u_{1,x}}{h(u, u_1, \alpha)} + \frac{u_{-1,x}}{h(u, u_{-1}, \alpha_{-1})} + \frac{\partial}{\partial u} \log (h(u, u_1, \alpha)h(u, u_{-1}, \alpha_{-1}))$$

where  $h(u, v, \alpha)$  is the symmetric biquadratic polynomial in  $u, v$  such that  $r(u) = h_v^2 - 2hh_{vv}$ .

In the Hamiltonian form it appeared earlier in [12]

$$u_x = \frac{2h}{u_1 - v} + h_v, \quad v_x = \frac{2h}{u - v_{-1}} - h_u, \quad h = h(u, v, \alpha)$$

---

[14] V.E. Adler, A.B. Shabat. On a class of Toda chains. *Theor. Math. Phys.* **111:3** (1997) 647–657.

and is closely related also to the elliptic Volterra lattice [15]

$$u_x = \frac{h(u_1, u, \alpha)}{u_1 - u_{-1}} - h_{u_1}(u_1, u, \alpha).$$

and Sklyanin lattice [16, 10].

---

[15] R.I. Yamilov. On classification of discrete evolution equations. *Usp. Mat. Nauk* **38:6** (1983) 155–156.

[16] E.K. Sklyanin. On some algebraic structures related to Yang-Baxter equation. *Funkts. analiz* **16:4** (1982) 27–34.